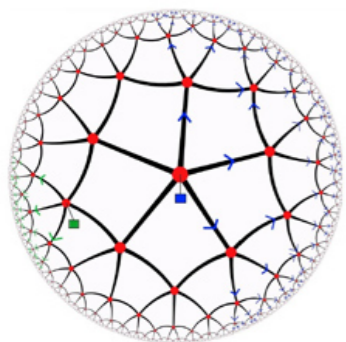


Holography and Entanglement

Matthew Headrick
Brandeis University

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ENTROPY AND AREA

Bekenstein-Hawking '74:

$$S = \frac{k_B c^3 \text{area}(\text{horizon})}{4G_N \hbar} = k_B \frac{\text{area}(\text{horizon})}{4l_P^2}$$

Planck length
($\approx 10^{-33}$ cm)

G_N → gravity

\hbar → quantum mechanics

k_B → statistical mechanics



Event Horizon Telescope '19



What are the “atoms” of the black hole?
Why is $S \propto \text{area}$?

ENTROPY AND AREA

If space has d dimensions, $G_N \hbar$ has units of area (L^{d-1})

Planck area: basic unit in quantum gravity, translates into unit of *entropy*

Generalizations of Bekenstein-Hawking:

De Sitter spacetime ([Gibbons-Hawking '77](#)):

$$S = \frac{\text{area}(\text{horizon})}{4G_N \hbar}$$

Holographic entropy bounds ([Bekenstein '81](#), [Bousso '99](#)): $S \leq \frac{\text{area}}{4G_N \hbar}$
for arbitrary closed surface in arbitrary spacetime

[Jacobson '95](#): area-entropy relation *implies* Einstein equation

How general is the area-entropy relation? What is its origin?

A clue: *Holographic entanglement entropy* ([Ryu-Takayanagi '06](#))

Vast (but also limited) generalization of Bekenstein-Hawking

To understand it, we first need to extend our notion of entropy...

ENTANGLEMENT ENTROPY

Classical mechanics:

definite state \rightarrow certain outcome for any measurement

Quantum mechanics:

definite state \rightarrow uncertain outcomes for some measurements

Example: $|\uparrow\rangle$

measurement of S_z definitely gives $+\frac{1}{2}\hbar$

measurement of S_x gives $+\frac{1}{2}\hbar$ or $-\frac{1}{2}\hbar$ with equal probability

When only certain kinds of measurements are allowed, a definite (pure) state will *effectively* be indefinite (mixed)

Suppose a system has two parts, but we can only measure one part

Spin singlet state: $|0, 0\rangle = \frac{1}{\sqrt{2}} (|\uparrow\rangle_A |\downarrow\rangle_B - |\downarrow\rangle_A |\uparrow\rangle_B)$ $S_{AB} = 0$

To see that this is a pure state (superposition, not mixture, of $|\uparrow\rangle|\downarrow\rangle$ and $|\downarrow\rangle|\uparrow\rangle$) requires access to both A and B

For an observer who only sees A , effective state is mixed:

$$\rho_A = \frac{1}{2} (|\uparrow\rangle\langle\uparrow| + |\downarrow\rangle\langle\downarrow|) \quad S_A = \ln 2$$

Indefiniteness from entanglement

ENTANGLEMENT ENTROPY

In general: if ρ is state of full system, then effective state for subsystem A is

$$\rho_A = \text{Tr}_{A^c} \rho$$

Entanglement entropy is defined as von Neumann entropy of subsystem:

$$S_A = -\text{Tr} \rho_A \ln \rho_A$$

(Any entropy can be viewed as due to entanglement with environment)

Entanglement entropies obey many important properties, such as:

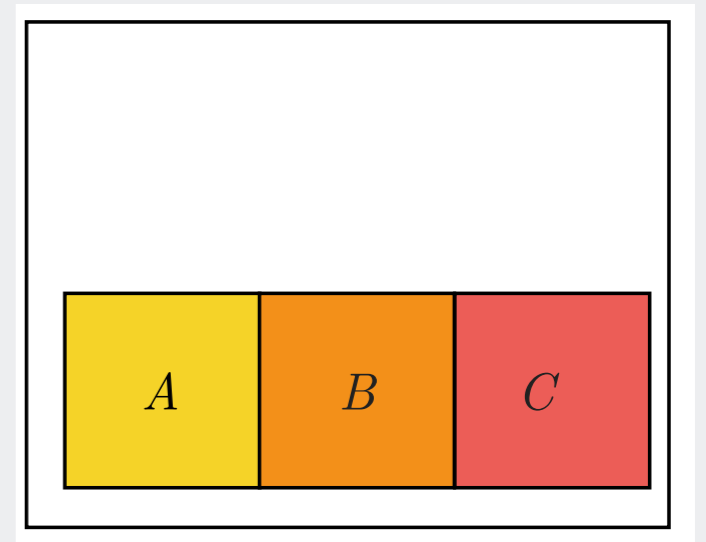
Subadditivity: $S_{AB} \leq S_A + S_B$

Mutual information: $I_{A:B} := S_A + S_B - S_{AB} \geq 0$

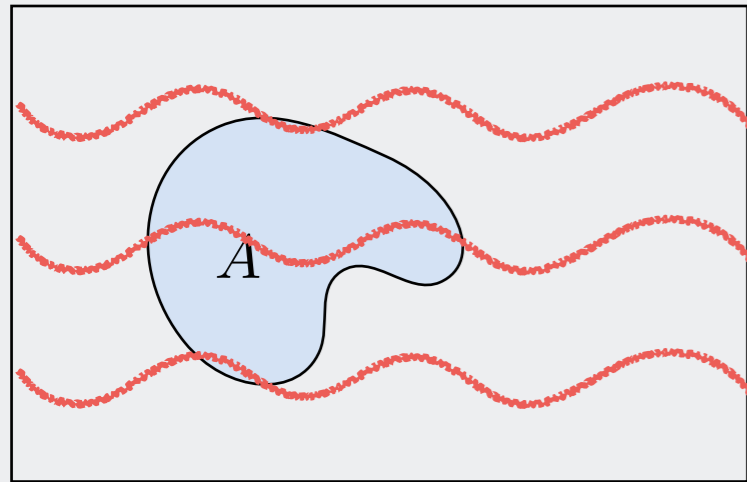
Strong subadditivity: $S_{AB} + S_{BC} \geq S_B + S_{ABC}$

(Lieb-Ruskai '73)

$$(I_{A:BC} \geq I_{A:B})$$



ENTANGLEMENT ENTROPY IN QFT



In quantum field theories (& many-body systems), spatial regions are highly entangled with each other

Consider microwave cavity

Even in vacuum, electromagnetic field fluctuates:
zero-point quantum fluctuations of modes

Each mode is distributed in space

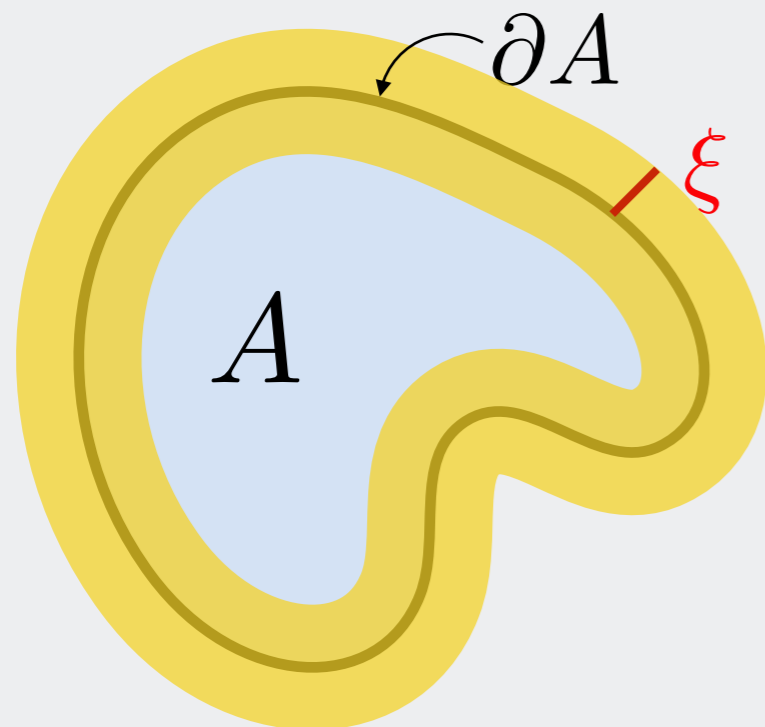
=> fluctuations are spatially correlated

=> any part A of cavity is entangled with rest

S_A is ultraviolet divergent due to entanglement of short-wavelength modes across ∂A

Massive (gapped) field: entanglement extends out to *correlation/Compton length*

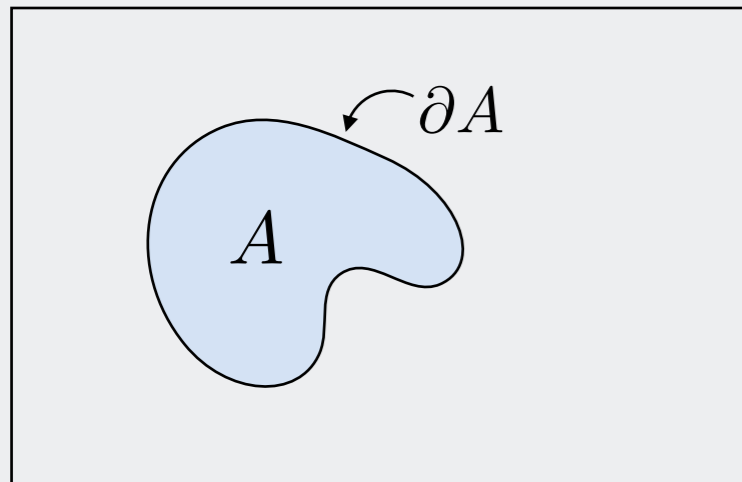
$$\xi \sim 1/m$$



$$S_A = \text{area}(\partial A) \left(\frac{1}{\epsilon^{d-1}} - \frac{1}{\xi^{d-1}} \right) + \dots$$

short-distance cut-off

ENTANGLEMENT ENTROPY IN QFT



In quantum field theories (& many-body systems), spatial regions are highly entangled with each other

S_A depends on:

- parameters of theory (including ϵ)
- state
- size and shape of region A

Contains a lot important physics

Examples:

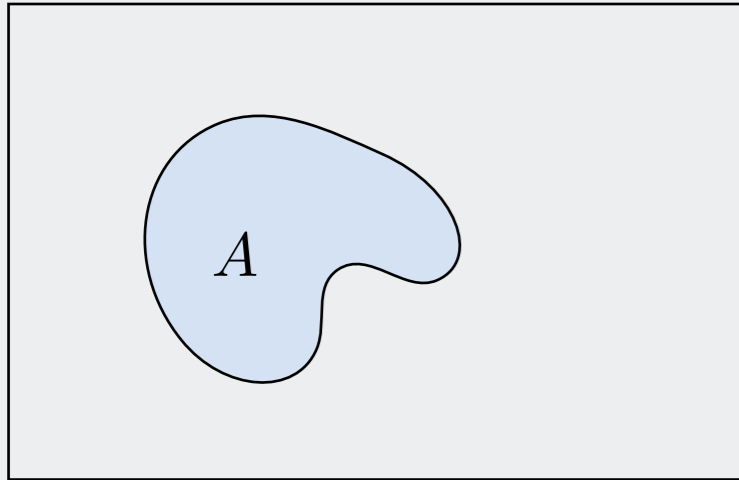
- Gapped theory in $d = 2$: $S_A = \frac{L}{\epsilon} - \frac{L}{\xi} - \gamma$
 - $\frac{L}{\epsilon}$: length(∂A)
 - $\frac{L}{\xi}$: correlation length
 - γ : topological entanglement entropy (Kitaev-Preskill '05; Levin-Wen '05)

- Critical (conformal) theory in $d = 1$: $S_A = \frac{c}{3} \ln \frac{L}{\epsilon}$
 - c : central charge
 - $\frac{L}{\epsilon}$:

A diagram of a horizontal line representing a 1D system. A blue shaded segment of length L is labeled 'A'. A small distance ϵ is indicated at the right end of the segment, representing a short-distance cutoff.
 - ϵ : short-distance cutoff

- At finite temperature, also usual extensive entropy: $s(T) \times \text{volume}(A)$
 - $s(T)$: thermal entropy density

ENTANGLEMENT ENTROPY IN QFT



Powerful probe of QFTs and many-body systems:

- quantum criticality
- topological order
- renormalization-group flows
- energy conditions
- many-body localization
- quenches
- much more...

However, usually very difficult to compute—even in free theories

Simplifies in certain theories with *many strongly-interacting* fields...

HOLOGRAPHIC DUALITIES

Consider a QFT with N interacting fields

for example $SU(n)$ Yang-Mills theory, $N \sim n^2$

When N is large, these fields may admit a *collective* description in terms of a small number of degrees of freedom

- classical (think of hydrodynamics)
- usually complicated

However, in certain cases, when the fields are very *strongly* interacting, it simplifies dramatically:

General relativity in $d + 1$ dimensions with cosmological constant $\Lambda < 0$

(plus some matter fields)

subject to certain boundary conditions: “universe in a box”

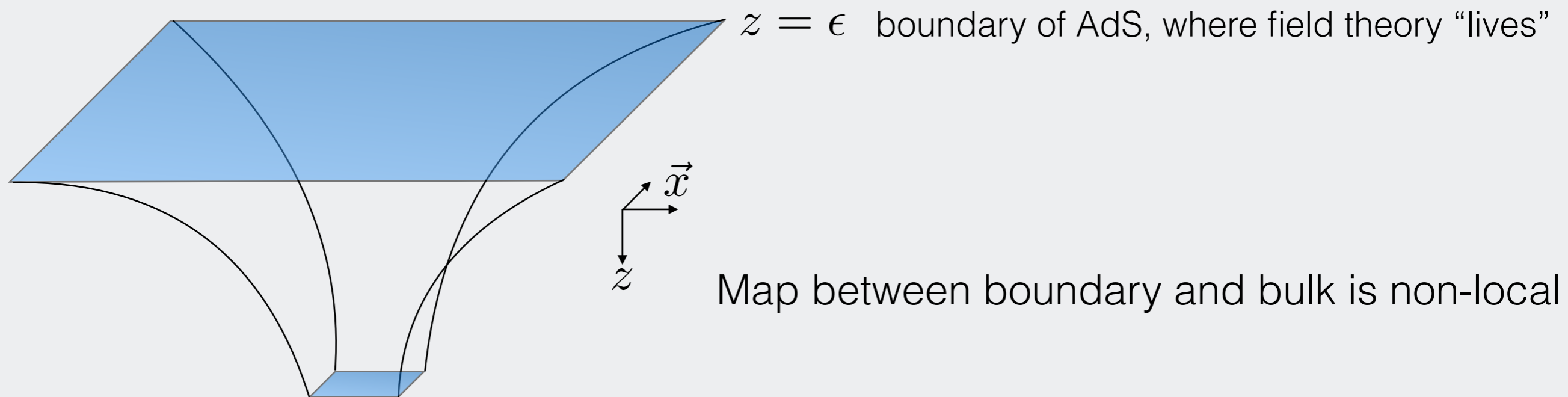
(Maldacena '97)

HOLOGRAPHIC DUALITIES

If QFT is conformal (scale-invariant), ground state is anti-de Sitter (AdS) spacetime:

$$ds^2 = \frac{R^2}{z^2} (-dt^2 + dz^2 + d\vec{x}^2)$$

AdS radius $\rightarrow R^2$
 d dimensions of QFT $\rightarrow d\vec{x}^2$
extra dimension $\rightarrow dz^2$

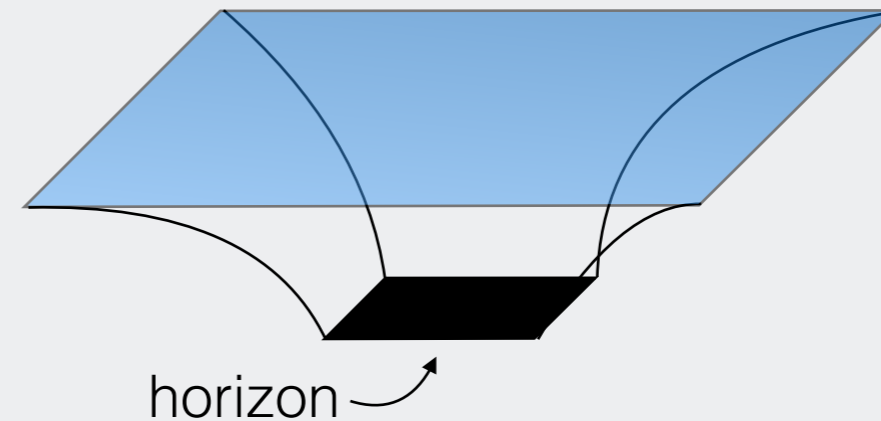


If QFT is gapped (massive), space ends on wall at $z_{\max} \sim \xi$ (correlation length)

Many specific examples known in various dimensions
(mostly supersymmetric, derived from string theory)

HOLOGRAPHIC DUALITIES

QFT	GR
N	AdS radius $\curvearrowright R^{d-1}$ Planck area $\curvearrowleft \frac{1}{G_N \hbar}$
thermodynamic limit $N \rightarrow \infty$	classical limit $\hbar \rightarrow 0$
statistical fluctuations	quantum fluctuations
collective modes	gravitational waves, etc.
deconfined plasma	black hole
$S \propto N$	$S = \frac{\text{area}(\text{horizon})}{4G_N \hbar}$



Holographic dualities are useful for computing *many* things in strongly interacting QFTs

Let's talk about entanglement entropies...

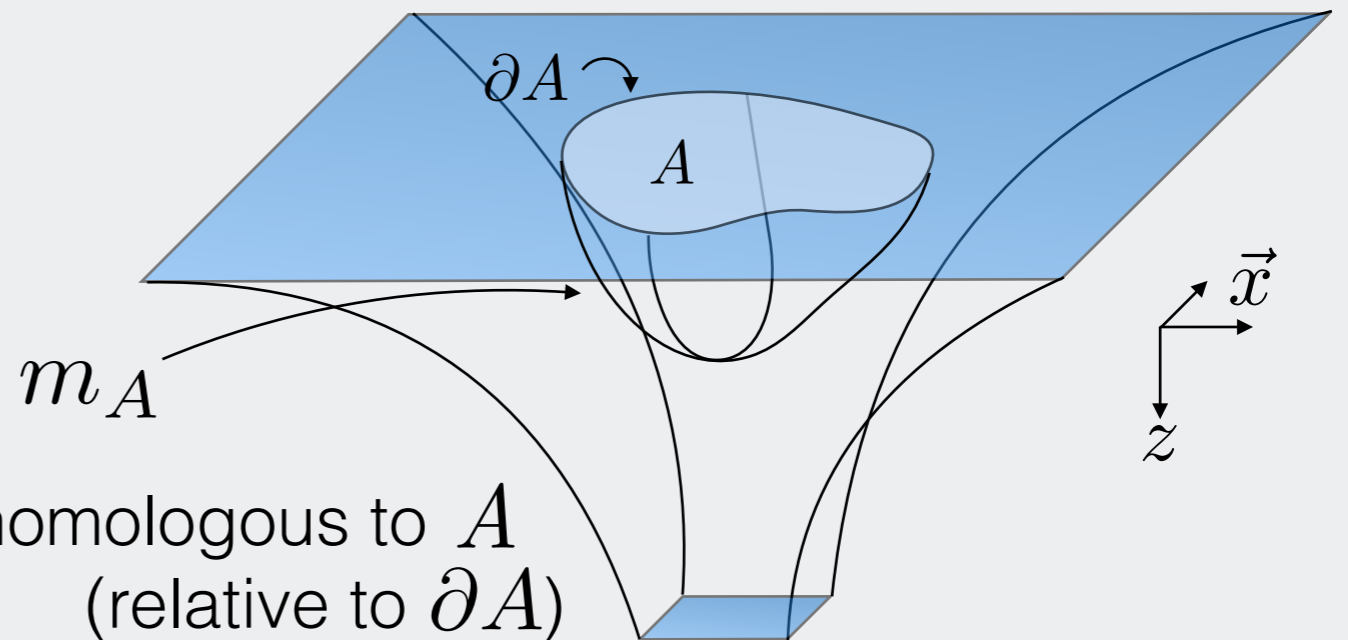
HOLOGRAPHIC ENTANGLEMENT ENTROPY

Ryu-Takayanagi '06:

$$S_A = \frac{\text{area}(m_A)}{4G_N \hbar}$$

m_A = minimal-area hypersurface homologous to A
(relative to ∂A)

hangs down in order to minimize area

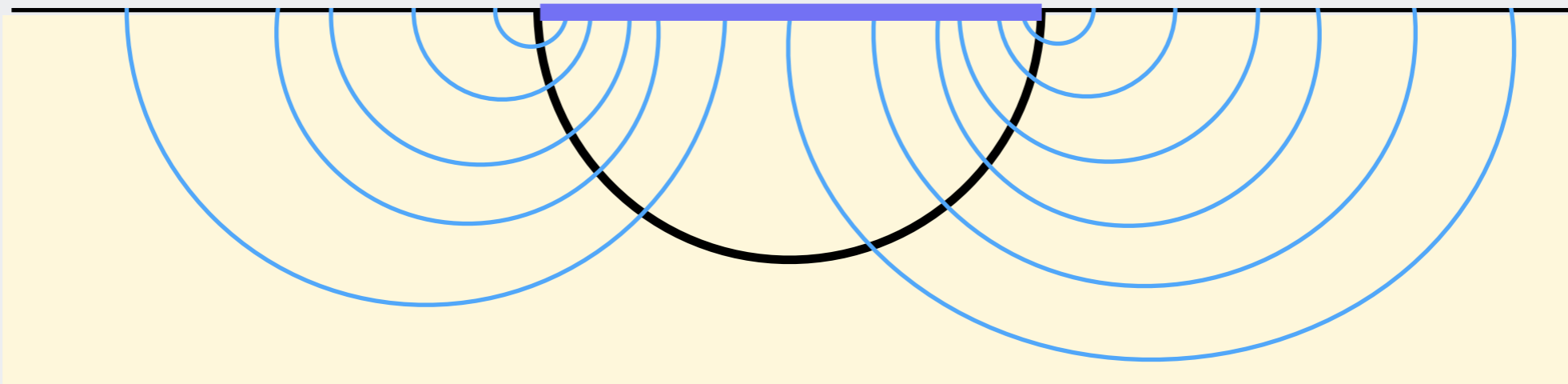


MH-Freedman '16: $S_A = \max \#$ “bit threads” connecting A to rest of boundary
(equivalence to minimal surface by Riemannian max flow-min cut theorem
Federer '74, ...)

Each bit thread has cross section of 4 Planck areas

Represents entangled pair of qubits between A and complement

A



HOLOGRAPHIC ENTANGLEMENT ENTROPY

Ryu-Takayanagi '06:

$$S_A = \frac{\text{area}(m_A)}{4G_N \hbar}$$

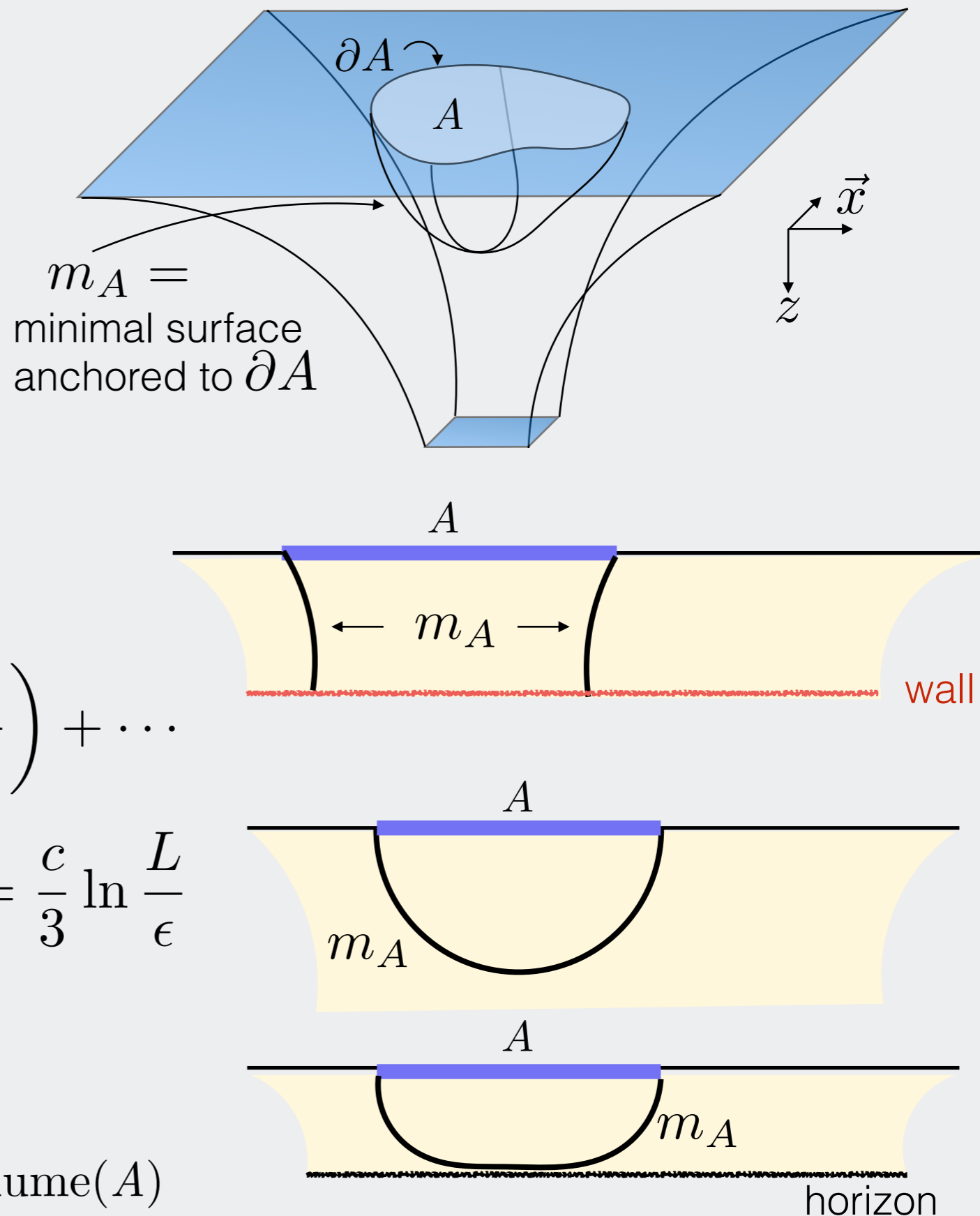
Geometrizes entanglement:

- Area-law UV divergence due to infinite area of m_A near boundary
- Gapped theory:
minimal surface extends to wall
(Klebanov, Kutasov, Murugan '07)

$$S_A = \text{area}(\partial A) \left(\frac{1}{\epsilon^{d-1}} - \frac{1}{\xi^{d-1}} \right) + \dots$$

- Conformal theory in $d = 1$: $S_A = \frac{c}{3} \ln \frac{L}{\epsilon}$

- Finite temperature:
minimal surface hugs horizon
=> extensive entropy $s(T)\text{volume}(A)$



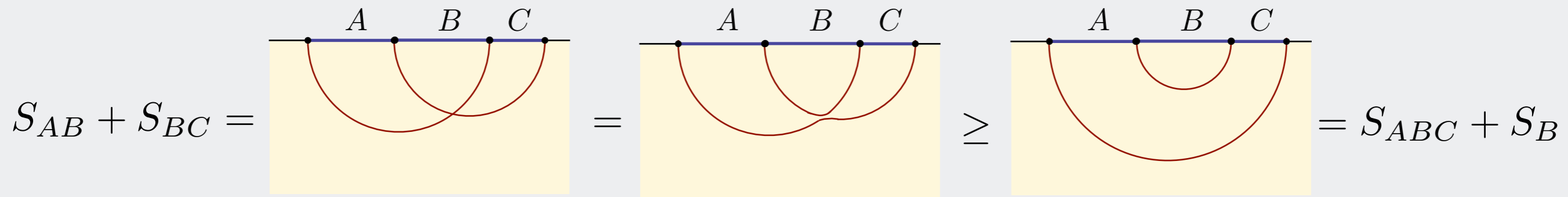
Democratizes Bekenstein-Hawking: not about horizons!

HOLOGRAPHIC ENTANGLEMENT ENTROPY

Quantum information theory is built into *classical spacetime geometry*

Example: Strong subadditivity

$$S_{AB} + S_{BC} \geq S_{ABC} + S_B$$



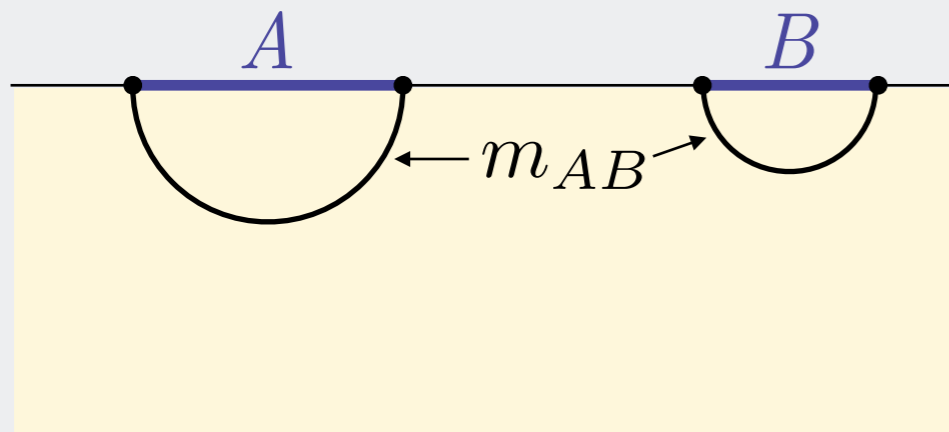
(MH-Takayanagi '07)

(Dual proof using bit threads [Freedman-MH '16](#))

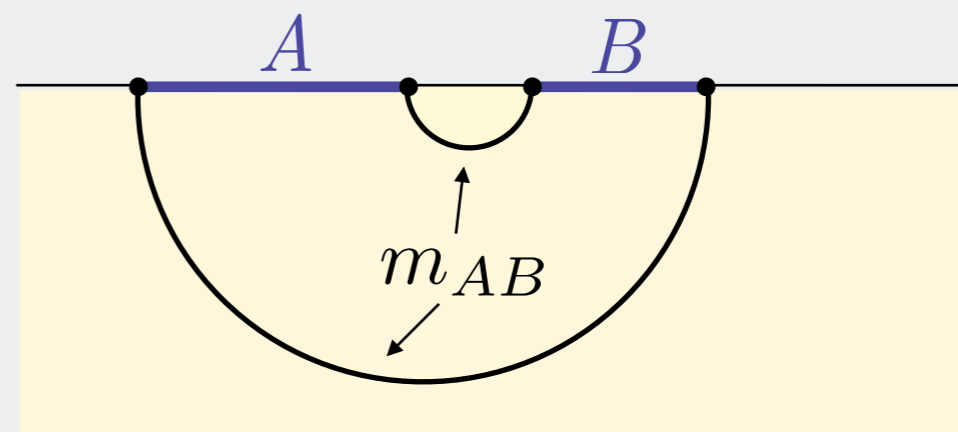
In fact, RT formula obeys *all* general properties of entanglement entropies
([Hayden-MH-Maloney '11](#); [MH '13](#))

HOLOGRAPHIC ENTANGLEMENT ENTROPY

Also has special properties, such as phase transitions (MH '10)



$$S_{AB} = S_A + S_B \quad \Rightarrow \quad I_{A:B} = 0$$



$$S_{AB} < S_A + S_B \quad \Rightarrow \quad I_{A:B} > 0$$

Monogamy of mutual information inequality (Hayden-MH-Maloney '11; MH '13):

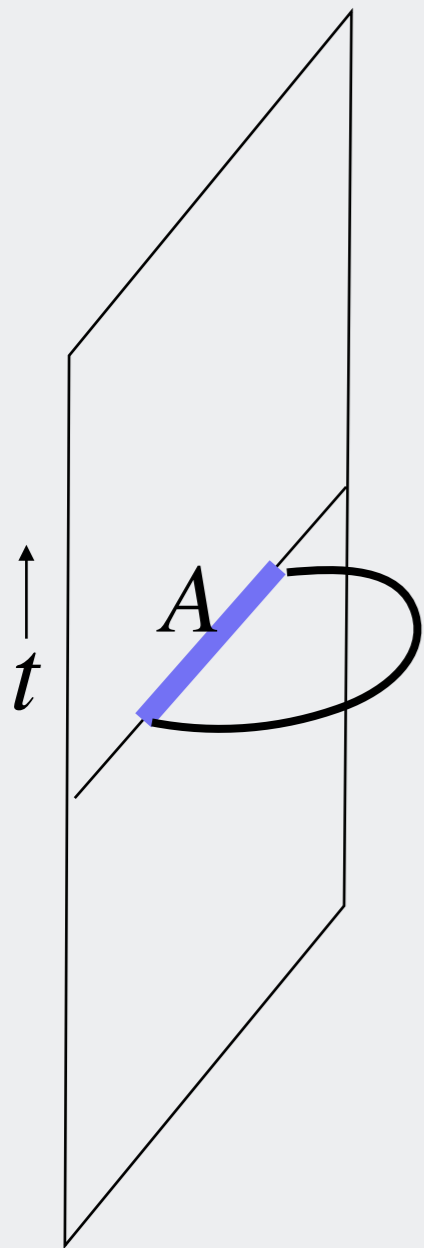
$$S_{AB} + S_{AC} + S_{BC} \geq S_A + S_B + S_C + S_{ABC}$$

$$I_{A:BC} \geq I_{A:B} + I_{A:C}$$

HOLOGRAPHIC ENTANGLEMENT ENTROPY

So far, we've ignored time

In *spacetime*, find minimal-area codimension-2 spacelike surface homologous to A
(Hubeny-Rangamani-Takayanagi '07)



Bulk spacetime geometry (Einstein equation) directly related to quantum information properties in field theory:

- Strong subadditivity (Wall '12)
- Causality (no faster-than-light signalling)
(MH-Hubeny-Lawrence-Rangamani '14)
- First law of entanglement:

$$\Delta S = \langle H_{\text{mod}} \rangle$$

To some extent, Einstein equation can be *derived* from properties of entanglement (Lashkari-McDermott-Van Raamsdonk '13)

HOLOGRAPHIC ENTANGLEMENT ENTROPY

Many other developments:

- Relation between bulk & boundary modular Hamiltonians and relative entropies
(Jafferis-Lewkowycz-Maldacena-Suh '15)
- Derivation of RT formula (MH '10, Lewkowycz-Maldacena '13)
- Tensor networks for modelling holography (Swingle '08)
- Bit threads & entanglement structures (Cui et al '18)
- Holography as quantum error-correcting code (Almheiri-Dong-Harlow '14)
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