

Discrete scale invariance in quantum spin chains - a proposed experiment.

Vaughan Jones,
Vanderbilt

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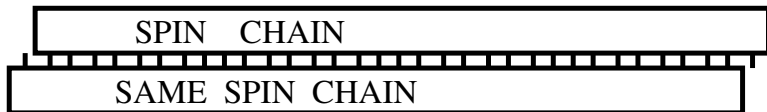
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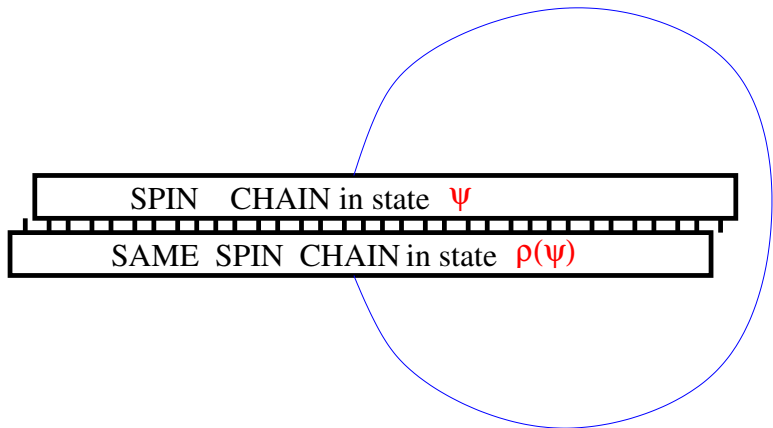
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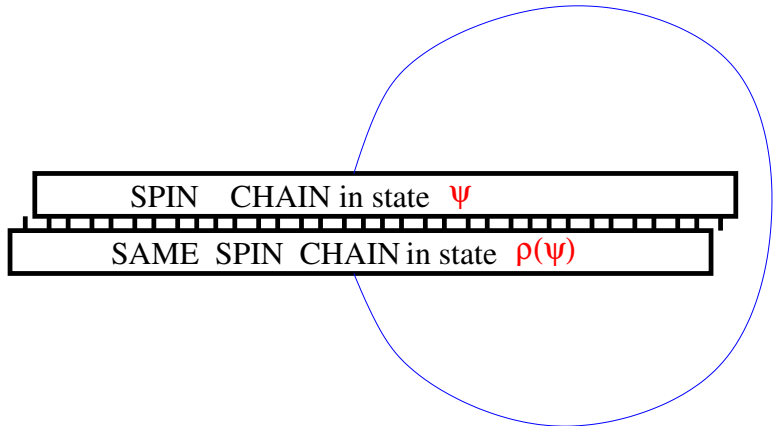


SPIN CHAIN in state ψ

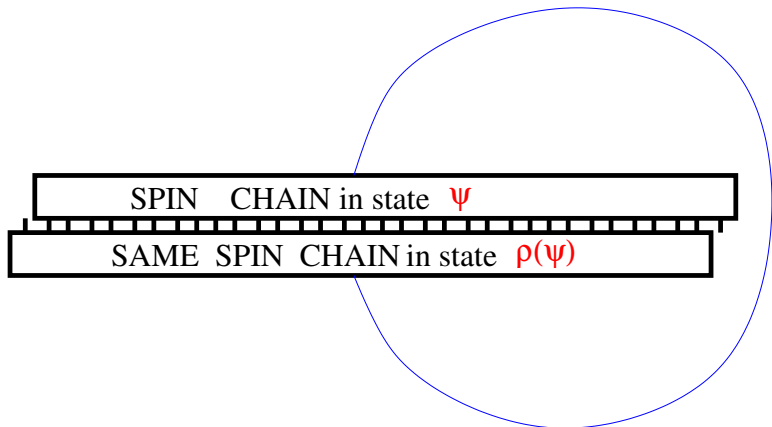
SAME SPIN CHAIN in state $\rho(\psi)$



Correlation almost 1 away from critical point.



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My "prediction": the correlation should *DROP* significantly at the critical point.

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For literally scale invariant (pure) states it is true-the result of a calculation I will give.

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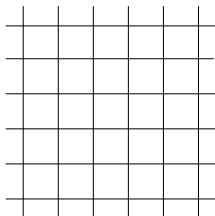
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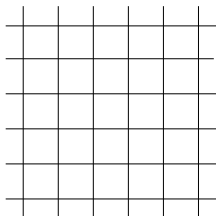
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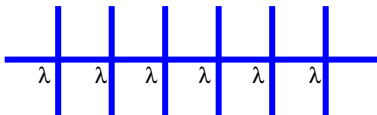
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One builds the lattice states up row by row with a matrix called the "Transfer matrix" $T(\lambda)$:



The rows and columns of $T(\lambda)$ are indexed by the possible spin states of a row of the lattice so that the terms in the product $T(\lambda)^n$ are enumerated by the spin states of an n -row lattice.

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Thus: "the transfer matrix determines the infinitesimal time evolution of the chain."

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This idea, that constrained quantum systems should be described by a relative tensor product, should not be restricted to one dimension.

Theorem

There exist bimodules ${}_M\mathcal{H}_M$ so that $(\overset{\circ}{\otimes})_M^n \mathcal{H}$ is finite dimensional and grows like $(2 \cos \pi/k)^n$ for $k = 3, 4, 5, \dots$.

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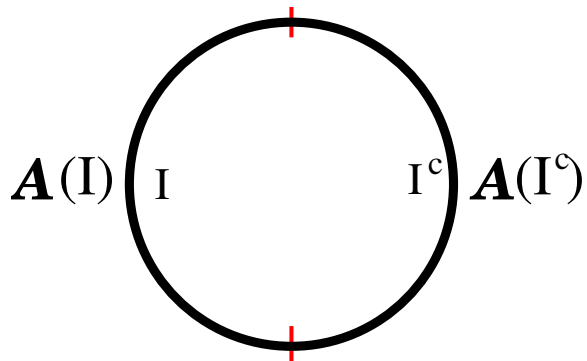
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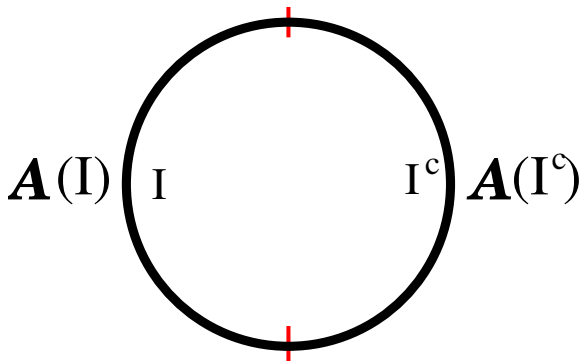
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.Wassermann's calculation of the Connes tensor product uses the KZ equation.

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The observables in $A(I)$ commute with those in $A(I^c)$ so we have a bimodule as required for our anyonic spin chain.

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My own approach has been inspired by the **diagrammatic, planar** nature of the structure of bimodules over a factor and the idea that quantum field theories should occur as the "scaling limit" of lattice models. And as we have seen, the bimodules in question provide anyonic quantum spin chains. This is not especially naive but faces huge mathematical difficulties.

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Thus one is led to create a Hilbert space associated to the circle by reversing block spin renormalisation and embedding the Hilbert space for any anyonic chain of length, say, 2^n inside one of length 2^{n+1} by **doubling all the spins**. For this one needs an elementary "spin doubling" operator



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A hallmark of CFT is the presence of representations of $Diff(S^1)$ on the Hilbert space. It was originally hoped that the Thompson group representations would tend as the lattice spacing tends to zero, to an action of $Diff(S^1)$.

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The calculation that yields this discontinuity is however very suggestive.

It involves **iterating a classical dynamical system (which can even be a rational transformation of the Riemann sphere) in the spectral parameter space of a transfer matrix!**

To approach the continuous limit we need to investigate how the ROTATION $\rho_{\frac{1}{2^n}}$ by $\frac{1}{2^n}$, which is an element of Thompson's group T , acts on states. In particular I want to calculate the coefficients

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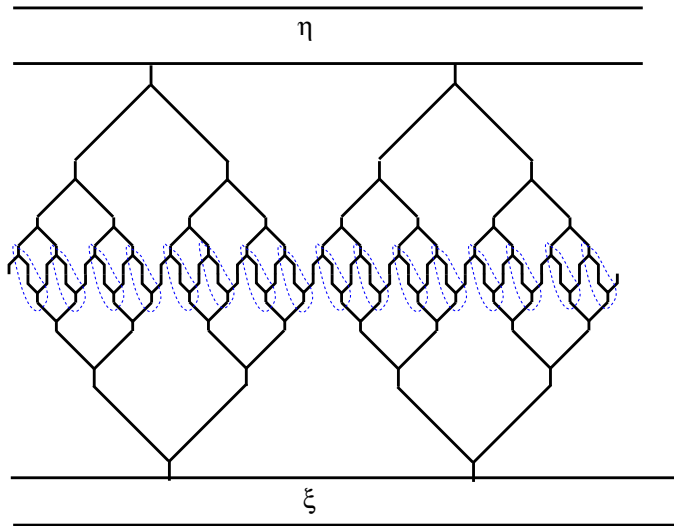
$$\langle \rho_{\frac{1}{2^n}}(\xi), \eta \rangle$$

Suppose that ξ and η are actually in some space $\otimes^{2^k} \mathcal{H}$. The following picture is $\langle \rho_{\frac{1}{2^{k+n+1}}} \xi, \eta \rangle$ which we illustrate here for $k = 1$ and $n = 3$.

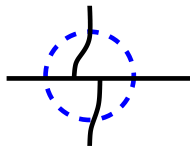
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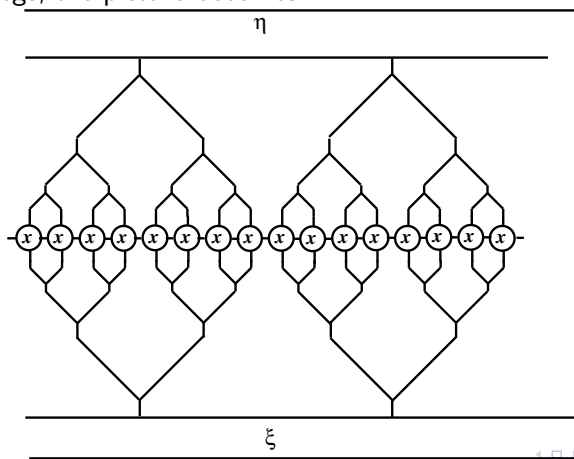


Now all the regions in the blue dotted circles can be deformed to look like



so if we call x the element inside the box with 4

legs, the picture becomes:



We recognise the *transfer matrix* $T_{2^{n+k}}(x)$!

Thus " The transfer matrix determines infinitesimal space translation". If we are in one dimension and time=space then we have recovered our previous mantra in a topsy turvy fashion!

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Note the resemblance between the calculation and the experiment.